

Solutions of Chi-square Quantile Differential Equation

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Abstract— The quantile function of probability distributions is often sought after because of their usefulness. The quantile function of some distributions cannot be easily obtained by inversion method and approximation is the only alternative way. Several ways of quantile approximation are available, of which quantile mechanics is one of such approach. This paper is focused on the use of quantile mechanics approach to obtain the quantile ordinary differential equation of the Chi-square distribution since the quantile function of the distribution does not have close form representations except at degrees of freedom equals to two. Power series, Adomian decomposition method (ADM) and differential transform method (DTM) was used to find the solution of the nonlinear Chi-square quantile differential equation at degrees of freedom equals to two. The approximate solutions converge to the closed (exact) solution. Furthermore, power series method was used to obtain the solutions for other degrees of freedom and series expansion was obtained for large degrees of freedom.

Index Terms— Chi-square, quantile function, differential equation, shape parameter, Adomian decomposition method, differential transform method.

I. INTRODUCTION

THE search for analytic expression of quantile functions has been a subject of intense research due to the importance of quantile functions. Several approximations are available in literature which can be categorized into four, namely: functional approximations, series expansions, numerical algorithms and closed form written in terms of a quantile function of another probability distribution which can also be refer to quantile normalization. In general, the notion of approximation of the quantile functions have been discussed extensively by [1-6]

The quantile function of the Chi-square is very important in statistical estimation [7-8]. Moreover the aim of the paper is to apply the use of quantile mechanics approach proposed by [9] to obtain a nonlinear second order ordinary differential equation which can be termed as “Chi-square Quantile Differential Equation (CQDE)” using a transformation of the probability density function (PDF) of the Chi-square distribution. This is a step towards the

quantile approximation of the distribution. This is because distributions with shape parameters require two steps towards effective quantile approximation [10]. Chi-square is an example of such distribution.

The solution of CQDE is the major contribution of the paper. This was done by the use of the power series, ADM and DTM for the case where the degrees of freedom is equal to two. This is to validate the methods for other cases and to create room for result comparison since the quantile function of the Chi-square distribution has closed form representation at that instance. The power series was used to obtain solutions for the other degrees of freedom.

The quantile mechanics as mentioned earlier is series expansion method of quantile approximation and has been applied for normal distribution [9], beta distribution [9], gamma [11], hyperbolic [12], exponential [13] and student’s t [14].

II. FORMULATION

The probability density function of the chi-square distribution and the cumulative distribution function are given by;

$$f(x) = \frac{1}{2^k \Gamma(k/2)} x^{k-1} e^{-x/2}, k > 0, x \in [0, +\infty) \quad (1)$$

$$F(x, k) = \frac{\gamma\left(\frac{k}{2}, \frac{x}{2}\right)}{\Gamma\left(\frac{k}{2}\right)} = P\left(\frac{k}{2}, \frac{x}{2}\right) \quad (2)$$

where $\gamma(.,.)$ = incomplete gamma function and $P(.,.)$ = regularized gamma function.

The quantile mechanics (QM) approach was used to obtain the second order nonlinear differential equation. QM is applied to distributions whose CDF is monotone increasing and absolutely continuous. Chi-square distribution is one of such distributions. That is;

$$Q(p) = F^{-1}(p) \quad (3)$$

Where the function $F^{-1}(p)$ is the compositional inverse of the CDF. Suppose the PDF $f(x)$ is known and the differentiation exists. The first order quantile equation is obtained from the differentiation of equation (3) to obtain;

$$Q'(p) = \frac{1}{F'(F^{-1}(p))} = \frac{1}{f(Q(p))} \quad (4)$$

Since the probability density function is the derivative of the cumulative distribution function. The solution to equation

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(4) is often cumbersome as noted by Ulrich and Watson [15]. This is due to the nonlinearity of terms introduced by the density function f . Some algebraic operations are required to find the solution of equation (4).

Moreover, equation (4) can be written as;

$$f(Q(p))Q'(p) = 1 \quad (5)$$

Applying the traditional product rule of differentiation to obtain;

$$Q''(p) = V(Q(p))(Q(p))^2 \quad (6)$$

Where the nonlinear term;

$$V(x) = -\frac{d}{dx}(\ln f(x)) \quad (7)$$

These were the results of [9].

It can be deduced that the further differentiation enables researchers to apply some known techniques to finding the solution of equation (6).

The reciprocal of the probability density function of the chi-square distribution is transformed as a function of the quantile function.

$$\frac{dQ(p)}{dp} = 2^{\frac{k}{2}}(\Gamma(k/2))Q(p)^{1-\frac{k}{2}}e^{\frac{Q(p)}{2}} \quad (8)$$

Differentiate again to obtain;

$$\frac{d^2Q(p)}{dp^2} = 2^{\frac{k}{2}}(\Gamma(k/2)) \left[\begin{aligned} & Q(p)^{1-\frac{k}{2}}e^{\frac{Q(p)}{2}} \frac{dQ(p)}{2dp} + \\ & \left(1-\frac{k}{2}\right)Q(p)^{\frac{k}{2}}e^{\frac{Q(p)}{2}} \frac{dQ(p)}{dp} \end{aligned} \right] \quad (9)$$

Factorization is carried out;

$$\frac{d^2Q(p)}{dp^2} = 2^{\frac{k}{2}}(\Gamma(k/2)) \left[\begin{aligned} & Q(p)^{1-\frac{k}{2}}e^{\frac{Q(p)}{2}} \frac{dQ(p)}{2dp} \\ & + \left(\frac{2-k}{2}\right) \frac{Q(p)^{1-\frac{k}{2}}}{Q(p)^{1-\frac{k}{2}}} Q(p)^{\frac{k}{2}}e^{\frac{Q(p)}{2}} \frac{dQ(p)}{dp} \end{aligned} \right] \quad (10)$$

$$\frac{d^2Q(p)}{dp^2} = \frac{1}{2} \left(\frac{dQ(p)}{dp} \right)^2 + \left(\frac{2-k}{2Q(p)} \right) \left(\frac{dQ(p)}{dp} \right)^2 \quad (11)$$

The second order nonlinear ordinary differential equations is given as;

$$\frac{d^2Q(p)}{dp^2} = \left(\frac{1}{2} + \frac{2-k}{2Q(p)} \right) \left(\frac{dQ(p)}{dp} \right)^2 \quad (12)$$

With the boundary conditions; $Q(0) = 0, Q'(0) = 1$.

III. POWER SERIES SOLUTION

The cumulative distribution function and its inverse (quantile function) of the chi-square distribution do not have closed form. However, an analytical formula is

available for the cumulative distribution function of the Chi-square distribution when the degrees of freedom $k = 2$.

The formula is given as;

$$F(x) = 2(1 - e^{-\frac{x}{2}}) \quad (13)$$

The quantile function $Q(p)$ can be obtained as;

$$p = 2(1 - e^{-\frac{Q(p)}{2}}) \Rightarrow \frac{p}{2} = 1 - e^{-\frac{Q(p)}{2}} \quad (14)$$

Taking logarithm;

$$-\frac{Q(p)}{2} = \ln\left(1 - \frac{p}{2}\right) \quad (15)$$

The quantile function is given by;

$$Q(p) = \ln\left(1 - \frac{p}{2}\right)^{-2} \quad (16)$$

The exact solution (equation (16)) is compared with the approximate solution (equation (12)), to compare the convergence of the approximate solution to the exact. This is to create an avenue for comparison between the exact and approximate values and consequently examine the validity of the quantile mechanics.

When $k = 2$, equation (12) becomes;

$$2 \frac{d^2Q(p)}{dp^2} - \left(\frac{dQ(p)}{dp} \right)^2 = 0 \quad (17)$$

Alternatively, the PDF of the chi-square distribution at $k = 2$ can be used. The PDF is given as;

$$f(x) = e^{-\frac{x}{2}} \quad (18)$$

Applying the Quantile Mechanics approach to obtain;

$$\frac{dQ(p)}{dp} = e^{\frac{Q(p)}{2}} \quad (19)$$

$$\frac{d^2Q(p)}{dp^2} = \frac{1}{2} e^{\frac{Q(p)}{2}} \frac{dQ(p)}{dp} \quad (20)$$

$$\frac{d^2Q(p)}{dp^2} = \frac{1}{2} \left(\frac{dQ(p)}{dp} \right)^2 \quad (21)$$

Equation (21) is the same with equation (17).

The general power series solution of equations (12) or (16) is given by;

$$Q(p) = c_0 + c_1p + c_2p^2 + c_3p^3 + c_4p^4 + c_5p^5 + c_6p^6 + \dots = \sum_{n=0}^{\infty} c_n p^n \quad (22)$$

Differentiate equation (22);

$$Q'(p) = c_1 + 2c_2p + 3c_3p^2 + 4c_4p^3 + 5c_5p^4 + 6c_6p^5 + \dots = \sum_{n=1}^{\infty} n c_n p^{n-1} \quad (23)$$

Differentiate equation (23);

$$Q''(p) = 2c_2 + 6c_3p + 12c_4p^2 + 20c_5p^3 + 30c_6p^4 + 42c_7p^5 + \dots = \sum_{n=2}^{\infty} n(n-1)c_n p^{n-2} \quad (24)$$

Substitute equations (24) and (23) into (22) and collect like

terms.

$$\begin{aligned} \text{constant: } 4c_2 - 1 = 0 &\Rightarrow c_2 = \frac{1}{4} \\ p: 12c_3 - 4c_2 = 0 &\Rightarrow c_3 = \frac{1}{12} \\ p^2: 24c_4 - 4c_2^2 - 6c_3 = 0 &\Rightarrow c_4 = \frac{1}{32} \\ p^3: 40c_5 - 8c_4 - 12c_2c_3 = 0 &\Rightarrow c_5 = \frac{1}{80} \\ p^4: 60c_6 - 10c_5 - 16c_2c_4 - 9c_3^2 = 0 &\Rightarrow c_6 = \frac{1}{192} \\ p^5: 84c_7 - 12c_6 - 20c_2c_5 - 24c_3c_4 = 0 \\ &\Rightarrow c_7 = \frac{1}{448} \\ p^6: 112c_8 - 14c_7 - 24c_2c_6 - 30c_3c_5 - 16c_4^2 = 0 \\ &\Rightarrow c_8 = \frac{1}{1024} \end{aligned} \quad (25)$$

The coefficients are substituted into equation (22) to obtain the power series solution.

The power series solution of equation (17) or (21) is given by;

$$\begin{aligned} Q(p) = p + \frac{1}{4}p^2 + \frac{1}{12}p^3 + \frac{1}{32}p^4 + \frac{1}{80}p^5 \\ + \frac{1}{192}p^6 + \frac{1}{448}p^7 + \frac{1}{1024}p^8 + \dots \end{aligned} \quad (26)$$

This is the approximate solution of equation (16),

IV. NUMERICAL RESULTS

Adomian Decomposition Methods (ADM) and Differential Transform Methods (DTM) are used to confirm the results of the power series. This is achieved by using the methods to solve equation (12) and compare with the exact value that is equation (11). The methods are semi-analytic in nature and have been applied extensively in numerical analysis, computational fluid mechanics, rigid bodies' analysis, elasticity, mathematical finance, risk analysis and so on. Details on the theories, modifications and applications of ADM and DTM can be found in [16], [17], [18], [19], [20], [21], [22], [23], [24].

Adomian Decomposition Method

Writing equation (21) in the integral form gives;

$$u(p) = p + \int_0^p \int_0^p \left\{ \frac{1}{2} \left(\frac{dQ}{dP} \right)^2 \right\} dPdP \quad (27)$$

By ADM, the infinite series solution is of the form;

$$u(p) = \sum_{n=0}^{\infty} u_n(p) \quad (28)$$

Now using (27) in (21), we have;

$$\sum_{n=0}^{\infty} u_n(p) = p + \int_0^p \int_0^p \left\{ \frac{1}{2} \left(\frac{dQ}{dP} \right)^2 \right\} dPdP \quad (29)$$

In view of (29), the zeroth order term can be written as;

$$\sum_{n=0}^{\infty} u_n(p) = p$$

(30) while other terms can be determined using the recurrence relations

$$\sum_{n=0}^{\infty} u_n(p) = \int_0^p \int_0^p \left\{ \frac{1}{2} \left(\frac{dQ}{dP} \right)^2 \right\} dPdP \quad (31)$$

The nonlinear terms in equation (31) can be represented as

$$B = \left(\frac{dQ}{dP} \right)^2 \quad (32)$$

and the Adomian polynomials are computed as follows:

$$\begin{aligned} B_0 &= \left(\frac{dQ_0}{dP} \right)^2, \quad B_1 = 2 \frac{dQ_0}{dP} \frac{dQ_1}{dP}, \\ B_2 &= 2 \frac{dQ_0}{dP} \frac{dQ_2}{dP} + \left(\frac{dQ_1}{dP} \right)^2, \\ B_3 &= 2 \frac{dQ_0}{dP} \frac{dQ_3}{dP} + 2 \frac{dQ_1}{dP} \frac{dQ_2}{dP} \end{aligned} \quad (33)$$

Substituting equation (32) in equation (31) yields

$$\sum_{n=0}^{\infty} u_n(p) = \int_0^p \int_0^p \left\{ \frac{1}{2} B_n \right\} dPdP \quad (34)$$

Solving equations (30) and (34) yields the solution equation (21).

The series solution of equation (21) using the ADM is;

$$\begin{aligned} Q(p) = p + \frac{1}{4}p^2 + \frac{1}{12}p^3 + \frac{1}{32}p^4 \\ + \frac{1}{80}p^5 + \frac{1}{320}p^6 + \frac{11}{6720}p^7 + \dots \end{aligned} \quad (35)$$

Differential Transform Method

To solve the initial value problems by DTM we first transform equations (21) as;

$$U(k+2) = \frac{1}{(k+2)!} \left[\frac{1}{2} \sum_{r=0}^k (r+1)(k-r+1) U(k+1)U(k-r+1) \right] \quad (36)$$

and the initial condition as

$$U(0) = 0, U(1) = 1 \quad (37)$$

Using equation (37) in (36) and resolving the change of variables give the solution of equation (21)

The series solution of equation (21) using the DTM is;

$$\begin{aligned} Q(p) = p + \frac{1}{4}p^2 + \frac{1}{12}p^3 + \frac{1}{32}p^4 + \frac{1}{80}p^5 \\ + \frac{1}{192}p^6 + \frac{1}{448}p^7 + \frac{1}{1024}p^8 + \dots \end{aligned} \quad (38)$$

The three methods converge favorably to the exact value as shown in **Table 1**. That is the series solutions of equations (21), (35) and (38).

V. EXTENSION TO DIFFERENT DEGREES OF FREEDOM

The power series method was used to obtain the series solutions of equation (12) for the different degrees of freedom up to 20. No comparison was made because of the absence of the closed form of the CDF and QF. The result of the degrees of freedom equals to two is included. The coefficients of the series are shown in **Table 2**.

The equations formed a series which can be used to predict p for any given degree of freedom k.

$$Q(p) \approx p + \frac{1}{4(k-1)} p^2, \quad k > 1 \quad (39)$$

For very large k,

$$Q(p) \approx p \quad (40)$$

representation of both CDF and QF at that instance. The approximate solutions converge to the exact solution. The procedure serves as a validation mode for other degrees of freedom. The series solutions for up to degrees of freedom equal to 20 and for large cases were included. The methods used are efficient in handling nonlinear ODE and is recommended for solving quantile differential equations of probability distributions and most importantly quantile differential equations.

VI. CONCLUDING REMARKS

In this paper, the power series method, ADM and DTM was used to obtain the approximate solutions of the Chi-square quantile differential equations at degrees of freedom equals to two. Chi-square distribution has closed form

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Table 1: Numerical results of power series, ADM and DTM

p	EXACT	Power series	ERROR	ADM	ERROR	DTM	ERROR
0.05	0.050635616	0.050635616	9.08995E-16	0.05063562	3.87319E-09	0.050635616	4.16334E-17
0.10	0.102586589	0.102586589	4.54567E-13	0.102586712	1.22847E-07	0.102586589	8.32667E-17
0.15	0.155923083	0.155923083	1.78941E-11	0.155924007	9.24203E-07	0.155923083	8.32667E-17
0.20	0.210721031	0.210721031	2.44224E-10	0.210724888	3.8563E-06	0.210721031	5.55112E-17
0.25	0.267062785	0.267062783	1.86586E-09	0.267074431	1.16453E-05	0.267062785	1.11022E-16
0.30	0.325037859	0.325037849	9.87864E-09	0.325066511	2.86521E-05	0.325037859	5.55112E-17
0.35	0.384743785	0.384743745	4.06171E-08	0.384804964	6.11791E-05	0.384743785	1.66533E-16
0.40	0.446287103	0.446286964	1.38819E-07	0.446404815	0.000117713	0.446287103	1.66533E-16
0.45	0.509784499	0.509784087	4.12071E-07	0.509993582	0.000209082	0.509784499	1.11022E-16
0.50	0.575364145	0.57536305	1.09472E-06	0.575712658	0.000348513	0.575364145	1.77636E-15
0.55	0.643167248	0.643164589	2.65907E-06	0.643718787	0.000551539	0.643167248	1.54321E-14
0.60	0.713349888	0.713343888	5.99966E-06	0.714185623	0.000835735	0.713349888	1.15241E-13
0.65	0.786085176	0.786072449	1.27268E-05	0.787305392	0.001220216	0.786085176	7.4929E-13
0.70	0.861565832	0.861540212	2.56202E-05	0.863290668	0.001724835	0.861565832	4.27069E-12
0.75	0.940007258	0.939957947	4.93115E-05	0.942376246	0.002368987	0.940007258	2.16595E-11
0.80	1.021651248	1.021559954	9.12932E-05	1.02482115	0.003169903	1.021651247	9.9291E-11
0.85	1.106770476	1.106607093	0.000163383	1.110910758	0.004140281	1.106770476	4.16627E-10
0.90	1.195674002	1.195390177	0.000283824	1.200959055	0.005285053	1.195674	1.61689E-09
0.95	1.288714033	1.288233762	0.00048027	1.295311034	0.006597001	1.288714027	5.85495E-09

Table 2: Coefficients of the power series solution for different degrees of freedom

k	c_1	c_2	c_3	c_4	c_5	c_6	c_7
1	1	1	$\frac{5}{6}$	$\frac{7}{48}$	$\frac{871}{1440}$	$\frac{7131}{23040}$	$\frac{46977941}{4838400}$
2	1	$\frac{1}{4}$	$\frac{1}{12}$	$\frac{1}{32}$	$\frac{1}{80}$	$\frac{1}{192}$	$\frac{1}{448}$
3	1	$\frac{1}{8}$	$\frac{1}{36}$	$\frac{109}{576}$	$\frac{3167}{230400}$	$\frac{165701}{16588800}$	$\frac{3316213}{1219276800}$
4	1	$\frac{1}{12}$	$\frac{1}{72}$	$\frac{1301}{8640}$	$\frac{2647}{259200}$	$\frac{413599}{261273600}$	$\frac{484086437}{87787929600}$
5	1	$\frac{1}{16}$	$\frac{1}{120}$	$\frac{481}{3840}$	$\frac{23}{3000}$	$\frac{54733}{36864000}$	$\frac{254245339}{46448640000}$
6	1	$\frac{1}{20}$	$\frac{1}{180}$	$\frac{5407}{50400}$	$\frac{239}{40320}$	$\frac{12683}{13608000}$	$\frac{1332597397}{266716800000}$
7	1	$\frac{1}{24}$	$\frac{1}{252}$	$\frac{1135}{12096}$	$\frac{35857}{7620480}$	$\frac{5735}{9144576}$	$\frac{516617447}{118294235136}$
8	1	$\frac{1}{28}$	$\frac{1}{336}$	$\frac{523}{18816}$	$\frac{11803}{13171200}$	$\frac{782767}{12170188800}$	
9	1	$\frac{1}{32}$	$\frac{1}{432}$	$\frac{10373}{138240}$	$\frac{18017}{5702400}$	$\frac{2854367}{8758886400}$	
10	1	$\frac{1}{36}$	$\frac{1}{540}$	$\frac{5839}{85536}$	$\frac{105263}{42768000}$	$\frac{539797}{7506369000}$	
11	1	$\frac{1}{40}$	$\frac{1}{660}$	$\frac{3301}{52800}$	$\frac{239491}{58080000}$	$\frac{12257407}{48787200000}$	
12	1	$\frac{1}{44}$	$\frac{1}{792}$	$\frac{52285}{906048}$	$\frac{830068087}{226040855040}$	$\frac{829051693}{4068735390720}$	
13	1	$\frac{1}{48}$	$\frac{1}{936}$	$\frac{33703}{628992}$	$\frac{130147}{76658400}$	$\frac{513382873}{313992806400}$	
14	1	$\frac{1}{52}$	$\frac{1}{1092}$	$\frac{28397}{567840}$	$\frac{237119}{158995200}$	$\frac{5262391}{52706908800}$	
15	1	$\frac{1}{56}$	$\frac{1}{1260}$	$\frac{827}{17640}$	$\frac{60419}{52479000}$	$\frac{3006253}{158696496000}$	
16	1	$\frac{1}{60}$	$\frac{1}{1440}$	$\frac{129617}{2937600}$	$\frac{248951}{211507200}$	$\frac{16702817}{241118208000}$	
17	1	$\frac{1}{64}$	$\frac{1}{1632}$	$\frac{8705}{208896}$	$\frac{29683}{28113920}$	$\frac{15186127}{259097886720}$	
18	1	$\frac{1}{68}$	$\frac{1}{1836}$	$\frac{187291}{4744224}$	$\frac{339179}{355816800}$	$\frac{12716789}{254053195200}$	
19	1	$\frac{1}{72}$	$\frac{1}{2052}$	$\frac{55409}{1477440}$	$\frac{2544787}{2947492800}$	$\frac{95609147}{1600741232640}$	
20	1	$\frac{1}{76}$	$\frac{1}{2280}$	$\frac{86647}{2425920}$	$\frac{1049101}{1334256000}$	$\frac{261139993}{6996838464000}$	

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